Self-Seeding of a Gain-Switched Integrated Dual-Laser Source for the Generation of Highly Wavelength-Tunable Picosecond Optical Pulses

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Abstract—The authors demonstrate the generation of nearly transform-limited optical pulses that are wavelength tunable over almost 50 nm. The wide tuning range is obtained by self-seeding a gain-switched source containing two Fabry—Perot lasers, and employing a widely tunable Bragg grating in the feedback loop. The generated pulses exhibit Side-mode suppression ratios of 50 dB above and across the full tuning range.

Index Terms—Optical fiber communications, optical pulse generation, self-seeding, semiconductor laser, wavelength tunable.

I. INTRODUCTION

THE USE OF wavelength tunability as a means of providing dynamic provisioning, in next-generation photonic systems, is currently a key area that is attracting much attention [1]. This interest could be attributed to the convergence of data networking with multiwavelength optical networking, a natural outcome driven by the impending needs of the Internet. In addition to this development, current trends and technology maturity favor the deployment of optical communication systems, operating at line rates 40 Gb/s and beyond, thereby making it more likely that return-to-zero (RZ) coding be used for data transmission, as it is easier to compensate for dispersion and nonlinear effects in the fiber by employing soliton-like propagation [2]. Taking into account these moves toward tunable optical systems employing RZ coding, it is obvious that the development of a wavelength tunable source of short optical pulses will be of paramount importance for future wavelength division multiplexed (WDM), optical time division multiplexed (OTDM), and hybrid WDM/OTDM optical communication systems [3].

One of the simplest and most reliable techniques available to generate wavelength-tunable picosecond optical pulses involves the self-seeding of a gain-switched Fabry–Perot (FP) laser, and many experimental schemes have been reported [4], [5]. Self-seeding entails the use of a wavelength selective external cavity to reinject a small fraction of the output light back into the gain-switched FP laser at only one longitudinal mode frequency. Provided that the optical signal reinjected into the laser arrives during the buildup of an optical pulse in the FP laser, then a single-moded output pulse is obtained. An important charac-

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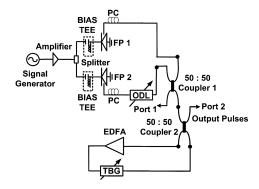


Fig. 1. Experimental setup for the generation of widely tunable optical pulses using the SSGS FP lasers.

teristic of these Self-Seeded Gain-Switched (SSGS) sources is the variation in the side-mode suppression ratio (SMSR) as the wavelength is tuned [6], as this may ultimately affect their usefulness in optical communication systems. In particular, recent work has demonstrated that, as the number of channels in a WDM system using SSGS pulse sources increases, the specifications on the required SMSR, due to cross-channel interference, may become very stringent [7]. It is thus likely that SMSRs in excess of 30 dB will be required for error-free operation of such systems.

Recent work in the development of wavelength tunable SSGS pulse sources has resulted in the generation of 90–130-ps optical pulses with SMSRs of around 32 dB that are tunable over 19–26 nm [8], [9]. The tunability of these systems was limited by factors such as the tunable range of the fiber Bragg grating (FBG) and the gain profile of the laser. In this letter, we show the generation of shorter pulses (~20 ps) that exhibit SMSRs greater than 50 dB over a tuning range approaching 50 nm. Our technique is based on the self-seeding of a gain-switched source containing two FP lasers, and the use of a widely tunable Bragg grating (TBG) filter. As we use two FP lasers with different gain curves, we can achieve a very large wavelength tuning range, and the high SMSR is essentially obtained by passing the self-seeded gain-switched pulses through the Bragg filter before the output.

II. EXPERIMENTAL SETUP

Fig. 1 illustrates the experimental configuration. It essentially consists of two gain switched FP laser diodes that are self-seeded using a single external cavity. The two FP lasers were chosen in such a way as to ensure that there was only

a very small overlap between their gain profiles, and finer placement of the gain profiles of the two lasers could be achieved by temperature controlling the diodes. The peak of the gain curve for FP 1 is at 1524 nm, while the peak of the gain curve for FP 2 is at 1561 nm. The incorporation of two gain-switched FP lasers enhances the wavelength tuning range that can be achieved using this self-seeding setup, as will be demonstrated below.

The FP lasers used were commercial 1.5-μm InGaAsP devices with threshold currents of approximately 26 mA, and mode spacings of 1.12 nm. Gain switching of both lasers was carried out by applying a dc-bias current of 17 mA and a 2.5-GHz sinusoidal modulation signal with a power of 29 dBm to both devices. The gain-switched output from both lasers were then coupled together before the composite signal was fed into an external loop cavity, which was used to self-seed both lasers. The external cavity contained a polarization controller (PC), a 3-dB coupler, a TBG (bandwidth: 0.23 nm, wavelength tuning range: 1460–1575 nm, rejection ratio: 40 dB, insertion loss: 5 dB) and an erbium-doped fiber amplifier (EDFA). An optical isolator in the EDFA ensures that light only propagates in one direction around the feedback loop, and the tunable filter eliminates unwanted amplified spontaneous emission from the EDFA in addition to selecting the laser mode to be seeded. The external cavity for self-seeding FP 2 also contains a tunable optical delay line (just at the output of the laser) to ensure that simultaneous self-seeding of FP 1 and FP 2 can be achieved. The EDFA in the external cavity is required to overcome the high losses obtained in the TBG (which has an insertion loss of approximately 5 dB across the tuning range and a very narrow linewidth compared to the spectral width (~63 GHz) of the modes from the gain-switched laser), and ensure that there is sufficient light reinjected into either laser to obtain suitable SMSRs on the output pulses.

To achieve optimum SSGS pulse generation, the Bragg grating was initially tuned to one of the longitudinal modes of the gain-switched FP-1 laser. The frequency of the modulation was then varied to ensure that the signal reinjected into the laser, from the external cavity, arrives as an optical pulse is building up in the laser. An operating frequency of 2.498 GHz was found to be suitable. The grating was then tuned to one of the longitudinal modes of FP 2 and, in this case, the tunable optical delay line was varied to ensure that the signal fed back into FP 2 arrives at the correct instant. The bias currents of FP 1 and FP 2 were then slightly changed to obtain the minimum pulsewidth. By subsequently tuning the grating across the gain curves of the lasers, we can achieve single-moded optical pulses over a very wide range of wavelengths. The output pulses, from the return arm of the second 50:50 fiber coupler (port 2) were characterized using a 50-GHz photodiode in conjunction with a 50-GHz digitizing oscilloscope, and a 0.05-nm resolution spectrum analyzer.

III. RESULTS AND DISCUSSION

The optical spectrum of the dual-wavelength signal from the gain-switched lasers without self-seeding is shown in Fig. 2(a). It can clearly be seen that, by combining the output of the gain

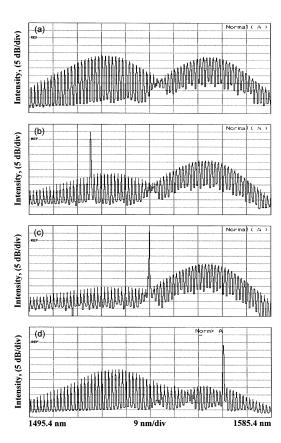


Fig. 2. Output optical spectra at port 1 of: (a) dual-wavelength signal. (b) Shortest wavelength that can be obtained: 1517.73 nm. (c) Central wavelength at 1540.4 nm. (d) Longest wavelength that can be obtained: 1566.64 nm.

switched lasers in the wavelength domain, the composite span of the gain profile that could be used for seeding has been greatly increased. The gain spectra of the two lasers overlap at approximately 16 dB down from the peak of their gain curves. This overlap corresponds to the maximum wavelength of FP 1, and the minimum wavelength of FP 2, for which we can achieve suitable SMSRs using the self-seeding configuration.

Different longitudinal modes of each FP laser were selectively excited when the seeding wavelength was tuned near the center of any desired mode. To obtain more continuous wavelength tuning with this setup, it is possible to use temperature tuning of the diodes in conjunction with tuning the Bragg grating. Fig. 2(b)–(d) shows in respective order the shortest, central, and longest wavelengths that could be seeded. The seeded spectra shown are the composite output (port 1) of the two SSGS lasers before passing through the optical filter and amplifier to be output to port 2.

By taking the output pulses at port 2, we thus pass the composite signal from the SSGS lasers through the external cavity before being outputted. The effect of this is to eliminate the signal from the unseeded laser, and greatly improve the SMSR of the generated optical pulses from the seeded laser [as shown below in Fig. 3(b) and (d)]. The output pulses, and their associated spectra, generated at two specific wavelengths (1524 and 1560 nm) are shown in Fig. 3. The deconvolved pulsewidth for the 1524-nm signal was 16 ps, while that of the 1560-nm signal was 18.5 ps. The associated spectral widths of these two

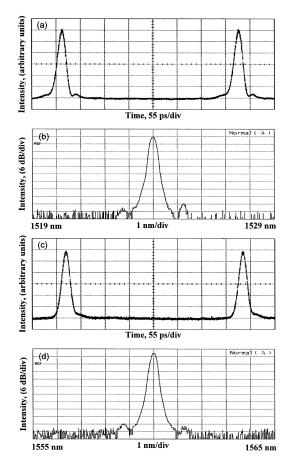


Fig. 3. (a) and (c) Output optical pulses and (b) and (d) associated spectra from port 2 for signal generation at 1524 nm [(a) and (b)] and 1560 nm [(c) and (d)].

signals were 27 and 26 GHz, respectively, while the SMSR of the generated pulses were 54 and 56 dB, respectively. The measured pulsewidth remained reasonably constant (16–20 ps) as the output pulses were tuned across the entire tuning range, with slight increases at the limits of tunability, and the time-bandwidth product of the generated pulses remains in the 0.43–0.49 range over the tuning range [which is close to that of a transform limited Gaussian pulses (0.44)]. The average output power of the optical pulses is approximately 1.8 mW. The main limitation on the wavelength tuning of the generated pulses was imposed by the gain bandwidth of the EDFA used in our experimental setup.

The dependence of the SMSR on the seeding wavelength was plotted and is shown in Fig. 4. It can be clearly seen that we were able to obtain an SMSR of 50 dB and above within a range of 48.91 nm (1517.73–1566.64 nm). As the seeding power was increased, due to higher pump powers from the EDFA, the achievable SMSR was enhanced and the possible tuning range became wider; however, pulse deformation and instabilities were observed. The results shown are taken at the optimum level of EDFA pump power (around 20 mW) to ensure maximum SMSR is achieved without pulse deformation and instabilities. Reduction of the EDFA pump power leads to a degradation in pulse SMSR, but this relationship is heavily dependent on the operating wavelength, and its position relative to the gain curve of the EDFA and FP laser being self-seeded. For example, at an operating wavelength of 1560 nm, where the

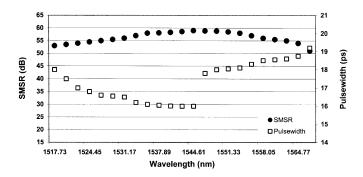


Fig. 4. SMSR (left-hand-side axis) and deconvolved pulsewidth (right-hand-side axis) against tunable range in the wavelength.

SMSR is 56 dB with 20-mW EDFA pump power, the SMSR is reduced to 49 dB and 44 dB at pump powers of 15 and 10 mW, respectively. Also presented in Fig. 4 is the pulsewidth variation as the wavelength is tuned. The point where the pulsewidth exhibits a sudden increase is the juncture when the seeded wavelength is moved from FP 1 to FP 2.

IV. CONCLUSION

The generation of widely tunable (\sim 50 nm) self-seeded gainswitched short optical pulses that exhibit very high SMSR (in the order of 50 dB) has been demonstrated. Such pulses (widely tunable and high SMSR) play a vital part in ensuring the optimal performance of high-speed WDM/OTDM optical communication networks [7]. Using an integrated dual laser source, it may be possible to develop a compact and highly stable SSGS wavelength tunable pulse source suitable for use in future high-speed optical networks.

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